

Chapter 15

Oscillation/Periodic Motion

Modified by P. Lam 8_2_2018

Learning Goals for Chapter 15

- To examine when natural oscillation occurs
 - To quantify and learn the details of a special type of oscillation called simple harmonic motion
 - To understand the concept of driven oscillations and resonance
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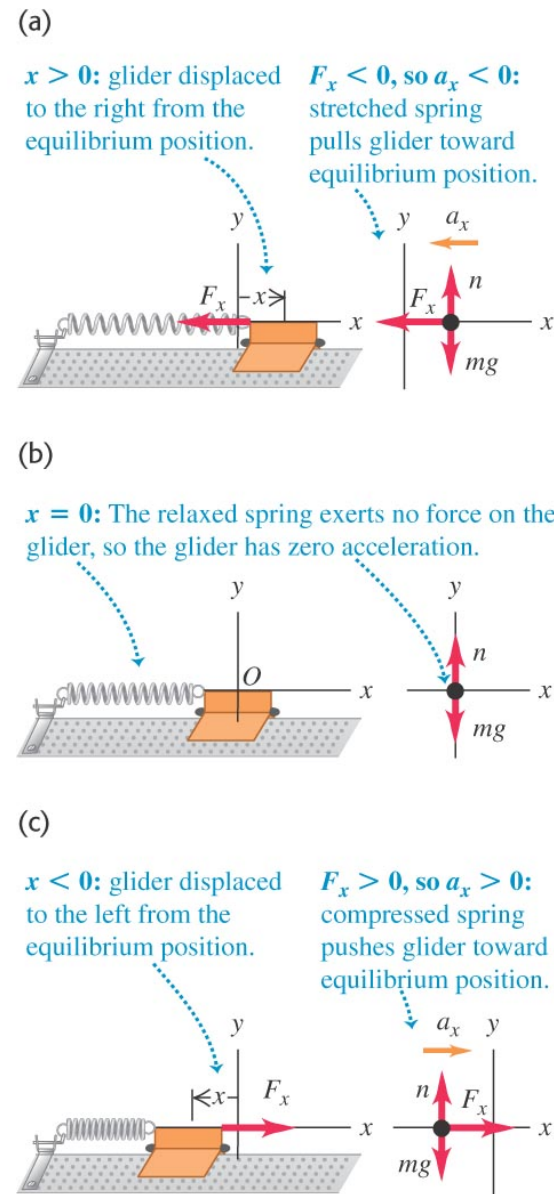
An example of natural oscillation

- A mass-spring system exhibits natural oscillation when the mass is displaced from its equilibrium position and then let go or give it push.

- Natural frequency

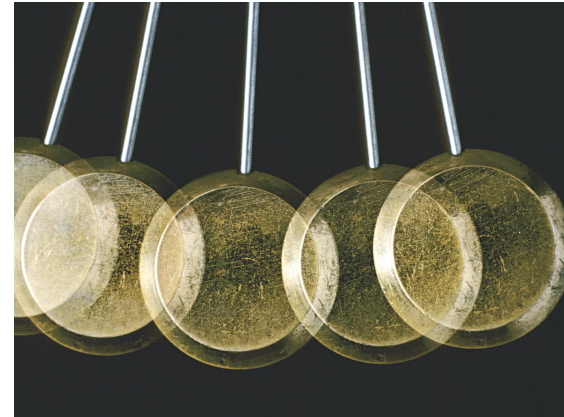
$$\omega = \sqrt{\frac{k}{m}} \text{ rad/s}$$

- The restoring force is provided by the stretched (or compressed) spring.



Introduction - when does natural oscillation occur?

- Example: A pendulum hanging straight down is in a stable equilibrium position.
- If you pull the pendulum to the side and then let go or give it a push, the pendulum will oscillate at its natural frequency,



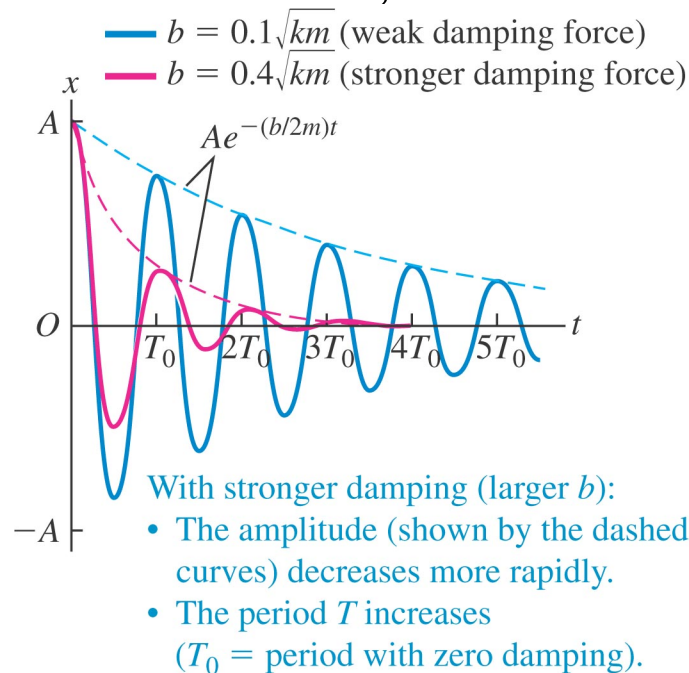
$$\omega = \sqrt{\frac{g}{\ell}} \text{ rad/s} \text{ (will derive later) if there is no damping forces.}$$

Friction, air-resistance, and other damping forces will cause the pendulum to oscillate at a slightly different frequency and to stop eventually.

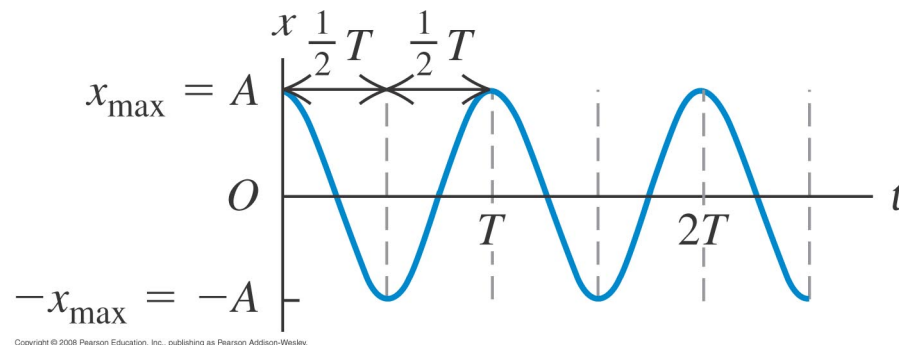
- Natural oscillation occurs whenever an object with mass (inertia) is displaced away from its ***stable equilibrium*** position. Note: the system provides a restoring force when the object is displaced from a stable equilibrium but not from an unstable equilibrium position.
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Different between oscillation and periodic motion

- The text equates periodic motion with oscillation.
- Technically, a damped oscillation is not a periodic function. The motion is periodic only when there is no damping force. In real situations, all natural oscillations are damped by friction, air-resistance, etc.



Damped oscillation

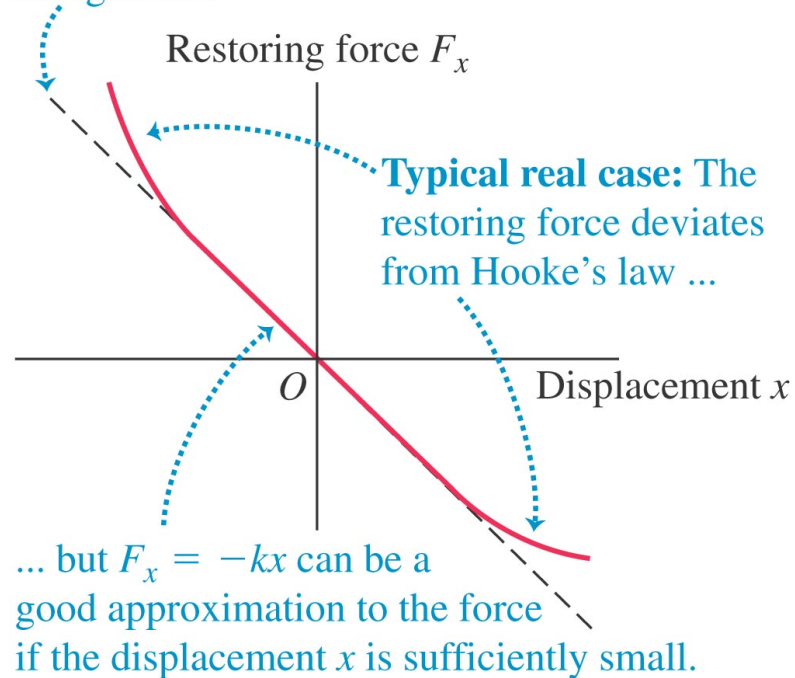


Undamped oscillation

A special type of oscillation - Simple harmonic motion

- An ideal spring responds to stretch and compression linearly, obeying Hooke's Law.

Ideal case: The restoring force obeys Hooke's law ($F_x = -kx$), so the graph of F_x versus x is a straight line.



Simple Harmonic motion occurs when there is no damping force and the restoring force is given by :

$$F = -kx \text{ (Hook's Law)}$$

For real springs, Hook's Law is a good approximation only when the displacement (x) is "small".

Mathematics of simple harmonic motion (SHM)

Newton's 2nd Law : $F = ma \Rightarrow -kx = m \frac{d^2 x}{dt^2}$

(assumed no damping force and restoring force obeys Hook's Law)

This equation is called a "differential equation" (D.E.).

That is an equation which relates the function, $x(t)$, to

its derivatives (in this case, $\frac{dx^2}{dt^2}$).

A differential equation is like a puzzle.

Try to think of a function, $x(t)$, which can satisfy the relationship stated by the differential equation.

In this case, what function whose 2nd derivative is proportional to the function with a negative sign?

Solution to the simple harmonic motion (SHM)

$$F = ma$$

$$-kx = m \frac{d^2 x}{dt^2}$$

Solution: $x(t) = A \cos \omega t + B \sin \omega t$

Substitute:

$$-k(A \cos \omega t + B \sin \omega t) = -m\omega^2 (A \cos \omega t + B \sin \omega t)$$

$\Rightarrow A$ & B can be any number

but

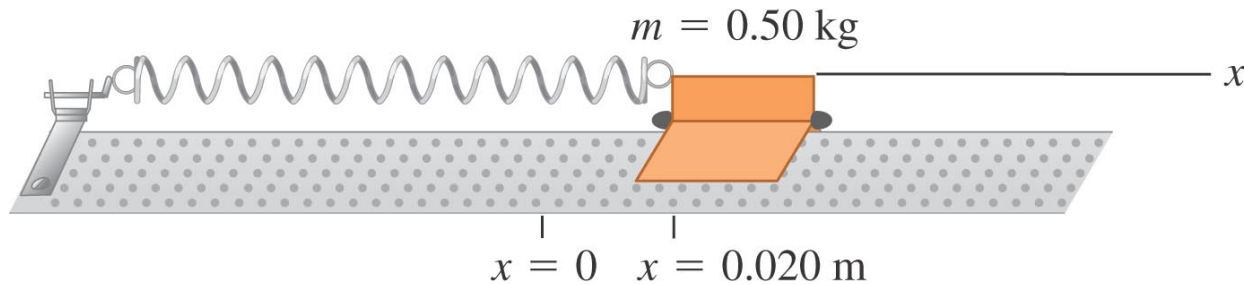
ω must be $\sqrt{\frac{k}{m}}$

ω is called the natural (angular) frequency of mass-spring system.

What are the physical meanings of A & B ?

Sample problems for SHM

(b)



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Given: Spring constant $k=800 \text{ N/m}$, $m=0.5\text{kg}$,

initial displacement $=+0.02\text{m}$, initial velocity $=0$

a) Find: frequency of oscillation in rad/s, period of oscillation in second.

b) Find $x(t)$ and sketch it.

c) Suppose the initial displacement $= 0.02\text{m}$, initial velocity $= +3\text{m/s}$,
find $x(t)$ and sketch it.

d) Suppose the initial displacement $= 0$, initial velocity $= +3\text{m/s}$,
find $x(t)$ and sketch it.

Solution to the simple harmonic motion (SHM) -cont

What does $x(t)=A \cos \omega t + B \sin \omega t$ represent?

What are the physical meanings of A & B?

Look at the situation at $t=0$

$x(t=0)=A \Rightarrow A$ is the initial displacement.

$$v(t)=\frac{dx}{dt} = -A\omega \sin \omega t + B\omega \cos \omega t$$

$$v(t=0) = B\omega \Rightarrow B = \frac{v_o}{\omega}$$

Note : As long as Hook's law is valid (i.e. the system exhibits simple harmonic motion),
the mass spring system will oscillate at the natural frequency no matter how large or small A or B.

Solution to the simple harmonic motion (SHM) -cont

Most general solution to $m \frac{d^2 x}{dt^2} = -kx$ is:

$$x(t) = A \cos \omega t + B \sin \omega t$$

$$= x_o \cos \omega t + \frac{v_o}{\omega} \sin \omega t$$

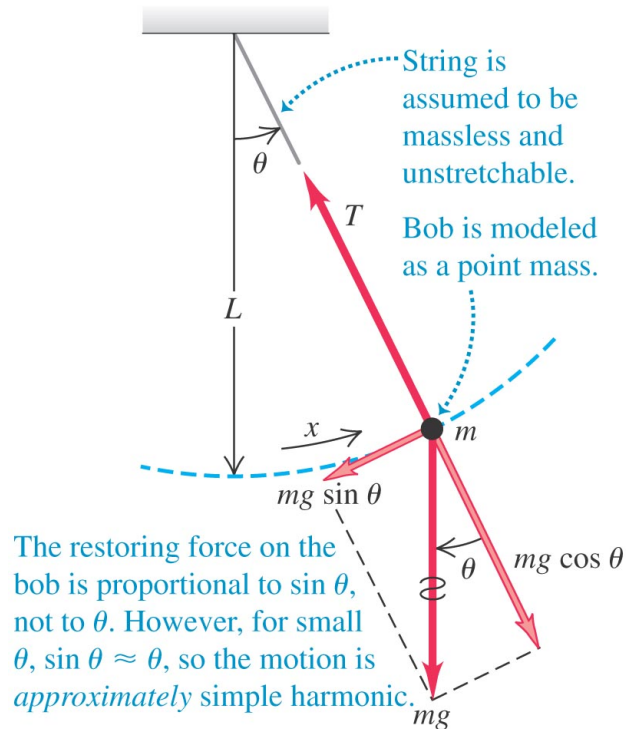
$$= \sqrt{x_o^2 + \frac{v_o^2}{\omega^2}} \cos(\omega t - \phi) \quad (\text{another way of writing the solution})$$

$$\tan \phi = \frac{v_o}{\omega x_o}$$

$$[\text{Note} : \cos(\omega t - \phi) = \cos \omega t \cos \phi + \sin \omega t \sin \phi]$$

Simple Pendulum

(b) An idealized simple pendulum



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$$\tau = I\alpha$$

$$-L(mg \sin \theta) = (mL^2) \frac{d^2\theta}{dt^2}$$

$$-\frac{g}{L} \sin \theta = \frac{d^2\theta}{dt^2}$$

$$\text{For "small" } \theta, \sin \theta \approx \theta \Rightarrow -\frac{g}{L} \theta = \frac{d^2\theta}{dt^2}$$

$$\text{compare with SHM: } -\frac{k}{m} x = \frac{d^2x}{dt^2}$$

\Rightarrow pendulum exhibits a SHM when θ is small (e.g. less than 1 radian)

\Rightarrow natural angular frequency for a simple pendulum is

$$\omega = \sqrt{\frac{g}{L}} \quad (\text{Note : does not depend on } m)$$

$$\Rightarrow \theta(t) = A \cos \omega t + B \sin \omega t$$

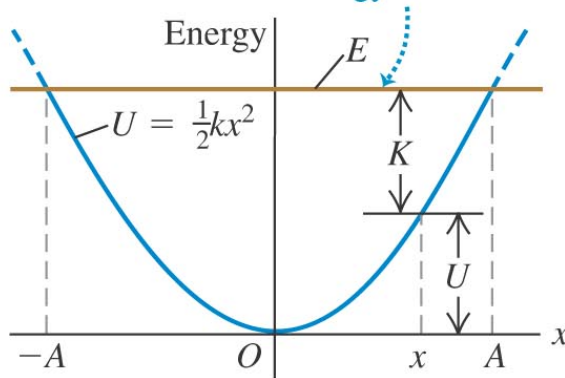
Q. Find $\theta(t)$ for a pendulum with mass=3kg, $L= 2.5\text{m}$, released from rest with an initial displacement is 0.5 rad.

Conservation of energy in SHM – graphical view

- Figure 13.15 shows the interconversion of kinetic and potential energy with an energy versus position graphic.

(a) The potential energy U and total mechanical energy E for a body in SHM as a function of displacement x

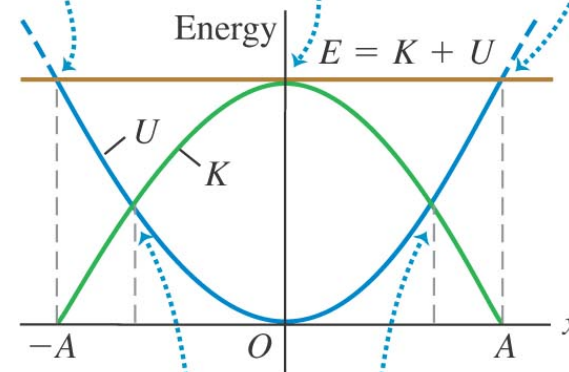
The total mechanical energy E is constant.



(b) The same graph as in (a), showing kinetic energy K as well

At $x = \pm A$ the energy is all potential; the kinetic energy is zero.

At $x = 0$ the energy is all kinetic; the potential energy is zero.



At these points the energy is half kinetic and half potential.

Conservation of energy in SHM

Given : $x(t) = x_o \cos \omega t + \frac{v_o}{\omega} \sin \omega t$ for SHM

Calculate the kinetic energy ($K = \frac{1}{2}mv^2$) and potential energy ($U = \frac{1}{2}kx^2$)

and show that $K+U$ does not change with time.

Try it yourselves first before looking at the answer below.

$$v(t) = \frac{dx}{dt} = -x_o \omega \sin \omega t + v_o \cos \omega t$$

$$K(t) = \frac{1}{2}m[v(t)]^2 = \frac{1}{2}m[(x_o \omega)^2 \sin^2 \omega t - 2x_o \omega v_o \sin \omega t \cos \omega t + (v_o)^2 \cos^2 \omega t]$$

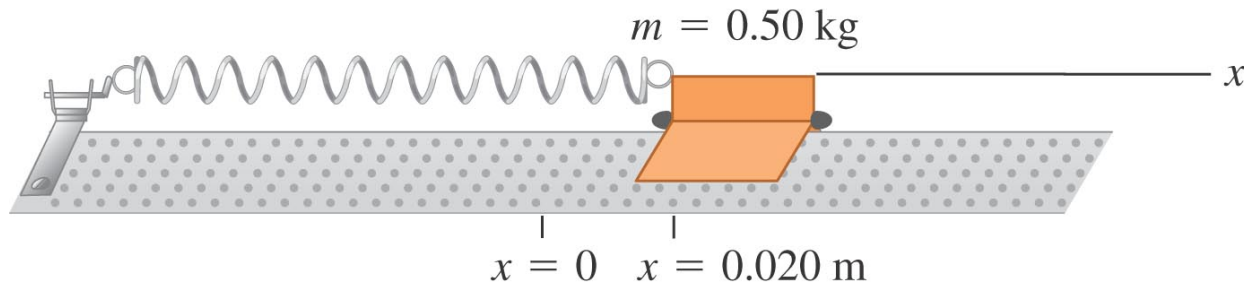
$$U(t) = \frac{1}{2}k[x(t)]^2 = \frac{1}{2}k\left[(x_o)^2 \cos^2 \omega t + 2x_o \frac{v_o}{\omega} \sin \omega t \cos \omega t + \left(\frac{v_o}{\omega}\right)^2 \sin^2 \omega t\right]$$

Note : $k = m\omega^2$

$$\Rightarrow K(t) + U(t) = \frac{1}{2}mv_o^2 + \frac{1}{2}kx_o^2 = K_o + U_o = \text{constant}$$

Energy Expression for SHM

(b)



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Given: Spring constant $k=800 \text{ N/m}$, $m=0.5\text{kg}$,
initial displacement $=+0.03\text{m}$, initial velocity $=2 \text{ m/s}$

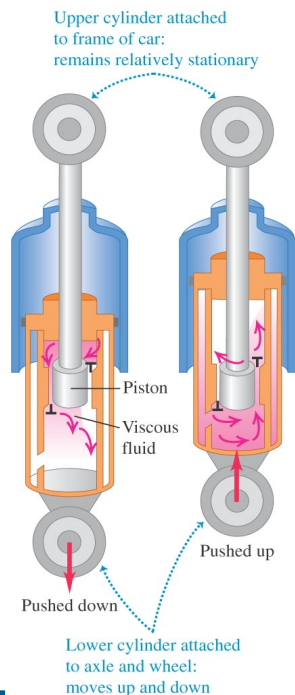
Find: The maximum displacement.

Find: The maximum velocity of the mass.

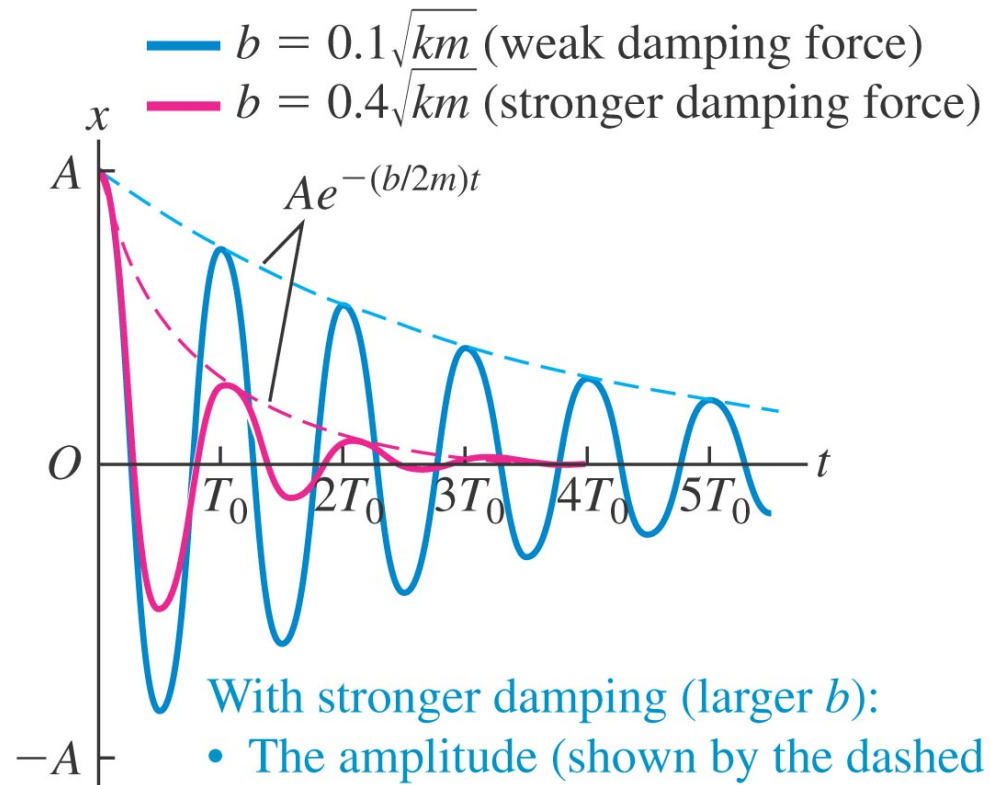
Find: The kinetic energy, potential energy and total energy at $t=2 \text{ s}$.

Damped oscillations

- Consider shock absorbers on your automobile. Without damping, hitting a pothole would set your car into SHM on the springs that support it.



$$m\ddot{x} = -kx - b\dot{x}; \quad b = \text{damping coefficient}$$



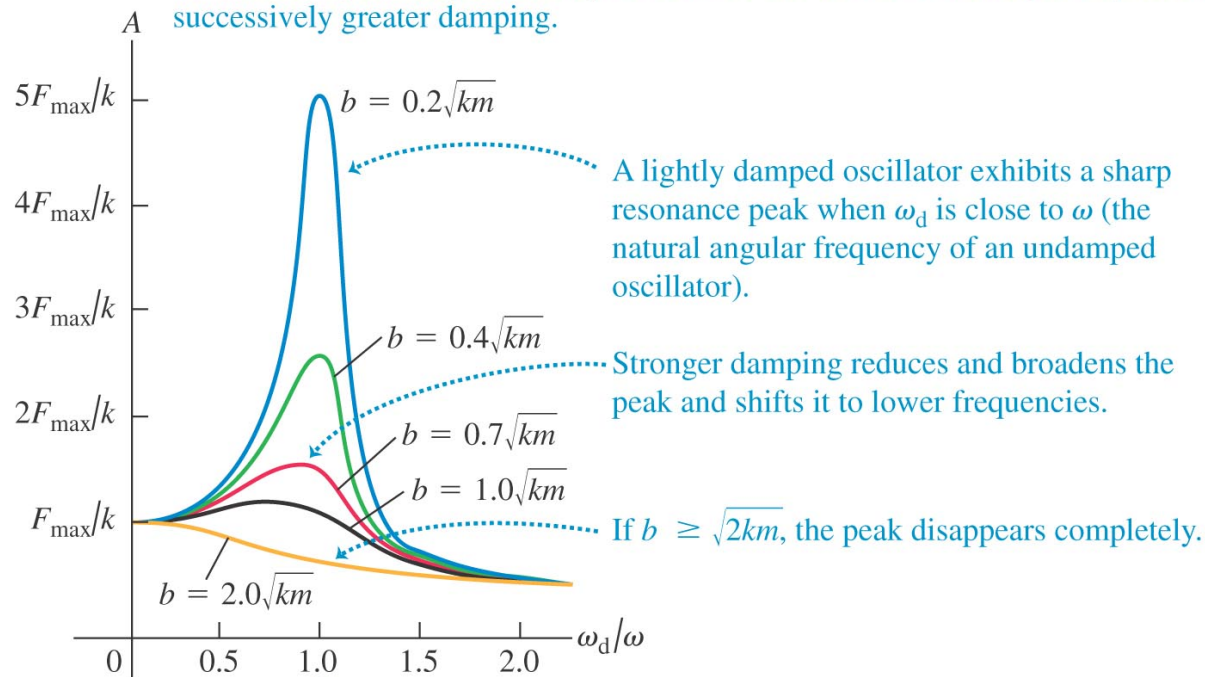
With stronger damping (larger b):

- The amplitude (shown by the dashed curves) decreases more rapidly.
- The period T increases ($T_0 =$ period with zero damping).

Forced (driven) oscillations and resonance

- A force applied “in synch” with a motion already in progress will resonate and add energy to the oscillation (refer to Figure 13.28).
- ***Resonance occurs when driven frequency=natural frequency***

Each curve shows the amplitude A for an oscillator subjected to a driving force at various angular frequencies ω_d . Successive curves from blue to gold represent successively greater damping.



Driving frequency ω_d equals natural angular frequency ω of an undamped oscillator.

The system has the largest amplitude of oscillation at resonance, see Figure to the left. Damping decreases the amplitude.

Mathematics of driven oscillator

$$m \frac{d^2 x}{dt^2} + kx = 0 \quad (\text{natural oscillation, natural freq.} = \omega_o = \sqrt{\frac{k}{m}})$$

$$m \frac{d^2 x}{dt^2} + kx = F_o \cos \omega t \quad (\text{driven oscillation, driven freq.} = \omega)$$

(assumed no damping force)

Solution: Let $x(t) = x_o \cos \omega t$ (system is forced to oscillate at driven freq.)

Substitute:

$$m(-\omega^2)x_o \cos \omega t + kx_o \cos \omega t = F_o \cos \omega t$$

$$\Rightarrow x_o = \frac{F_o}{k - m\omega^2} = \frac{F_o/m}{k/m - \omega^2} = \frac{F_o/m}{\omega_o^2 - \omega^2}$$

$\Rightarrow \infty$ amplitude when $\omega = \omega_o$

with damping force, amplitude is largest when $\omega = \omega_o$,
but not infinite.

Forced (driven) oscillations and resonance II

- The Tacoma Narrows Bridge suffered spectacular structural failure after absorbing too much resonant energy (refer to Figure 13.29).

